

Response function definitions:

$$C_V = \left(\frac{\partial E}{\partial T}\right)_V \quad C_P \equiv \left(\frac{\partial H}{\partial T}\right)_P \quad \kappa_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_T \quad \alpha_P = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P$$

Gases:

$$S(E, V, N) = k_B N \ln \left[ (E/N)^{3/2} (V/N) \right] + \frac{3}{2} k_B N \left( \frac{5}{3} + \ln \left( \frac{4\pi m}{3h^2} \right) \right)$$

$$\mu(T, P) = \mu^0(T) + k_B T \ln(yP) \quad \mu^0(T) \equiv -k_B T \ln \left( \frac{k_B T}{\Lambda(T)^3} \right) \quad \Lambda(T) \equiv \left( \frac{h^2}{2\pi m k_B T} \right)^{1/2}$$

$$\mu(T, P) = \mu^0(T) + k_B T \ln(yf) \quad \ln \frac{f(T, P')}{P'} = \int_0^{P'} \left( \frac{v(T, P)}{k_B T} - \frac{1}{P} \right) dP = \int_0^{P'} \frac{Z(T, P) - 1}{P} dP$$

Phase equilibrium:

$$\left(\frac{dP}{dT}\right)_{\text{phase boundary}} = \frac{\Delta h}{T \Delta v} \quad \ln P^{\text{vap}}(T) = c_1 - \frac{c_2}{c_3 + T}$$

Stability criteria:

$$\frac{\partial^2 F}{\partial X^2} > 0 \quad \frac{\partial^2 F}{\partial y^2} < 0$$

Solutions:

$$\mu_i(T, P, \{x\}) = \mu_i^*(T, P) + k_B T \ln(\gamma_i x_i)$$

$$y_i P = x_i P_i^{\text{vap}} \quad T_b' \approx T_b \left( 1 + \frac{k_B T_b}{\Delta h_{\text{vap}}} x_{\text{solute}} \right) \quad \Pi = \frac{k_B T x_{\text{solute}}}{v} \approx \frac{k_B T c_{\text{solute}}}{\mathcal{M}_{\text{solute}}}$$

$$-SdT + VdP = \sum_i N_i d\mu_i$$

$$\bar{X}_i = \left( \frac{\partial X}{\partial N_i} \right)_{T, P, N_{j \neq i}} \quad d\mu_i = -\bar{S}_i dT + \bar{V}_i dP + \sum_{k=1}^{C-1} \left( \frac{\partial \mu_i}{\partial x_k} \right)_{T, P, x_{j \neq k}} dx_k$$

$$-[x_1 \Delta S_1 + x_2 \Delta S_2] dT + [x_1 \Delta V_1 + x_2 \Delta V_2] dP + \left[ x_1 - x_2 \frac{y_1}{y_2} \right] \left( \frac{\partial \mu_1^G}{\partial y_1} \right)_{T, P} dy_1 = 0$$

Solids:

$$\frac{c_V^E}{3k_B} = \left( \frac{\Theta_E}{T} \right)^2 \frac{e^{\frac{\Theta_E}{T}}}{\left( e^{\frac{\Theta_E}{T}} - 1 \right)^2}, \quad \Theta_E \equiv \frac{h\nu}{k_B} \quad c_V^D(T, V) = 9k_B \left( \frac{T}{\Theta_D} \right)^3 \int_0^{\frac{\Theta_D}{T}} \frac{x^4 e^x}{(e^x - 1)^2} dx, \quad \Theta_D \equiv \frac{h\nu_m}{k_B}$$

Independent particles in the canonical ensemble:

$$Q = q^N \quad \text{or} \quad Q = \frac{q^N}{N!} \quad \text{where} \quad q \equiv \sum_x e^{-\beta \epsilon(x)}$$

Canonical fluctuations:

$$\wp(E) = \frac{\Omega(E, V, N) e^{-\beta E}}{Q(T, V, N)} \quad \langle E^v \rangle = \sum_E E^v \wp(E) = \frac{(-1)^v \partial^v Q}{Q \partial \beta^v} \quad \sigma_E^2 = C_V k_B T^2$$

Classical systems in the canonical ensemble:

$$\begin{aligned} \wp(p_x) &= (2\pi m k_B T)^{-\frac{1}{2}} \exp\left(-\frac{p_x^2}{2m k_B T}\right) & \langle K \rangle &= \frac{3}{2} N k_B T \\ \wp(v) &= 4\pi \left(\frac{m}{2\pi k_B T}\right)^{\frac{3}{2}} v^2 \exp\left(-\frac{m v^2}{2k_B T}\right) & \langle v \rangle &= \int v \wp(v) dv = \left(\frac{8k_B T}{\pi m}\right)^{\frac{1}{2}} \\ P &= \frac{N k_B T}{V} + \frac{1}{3V} \left\langle \sum \mathbf{f}_i \cdot \mathbf{r}_i \right\rangle \end{aligned}$$

Gibbs entropy formula:

$$S = -k_B \sum_m \wp_m \ln \wp_m$$

Chemical equilibrium:

$$\prod_i P_i^{v_i} = K_P(T) \quad \text{where} \quad K_P(T) \equiv \exp\left[-\frac{\sum_i v_i \mu_i^0}{k_B T}\right] \quad \frac{d \ln K_P}{d(1/T)} = -\frac{\Delta h^0}{k_B} \quad \text{where} \quad \Delta h^0 \equiv \sum_i v_i h_i^0$$

$$\prod_i \rho_i^{v_i} = K_C(T) \quad \text{where} \quad K_C(T) \equiv \exp\left[-\frac{\sum_i v_i \mu_i^0}{k_B T}\right] (k_B T)^{-\sum_i v_i}$$

$$\prod_i x_i^{v_i} = K_x(T, P) \quad \text{where} \quad K_x(T, P) \equiv \exp\left[-\frac{\sum_i v_i \mu_i^*}{k_B T}\right]$$

$$\prod_i \rho_i^{v_i} = K_C(T, P) \quad \text{where} \quad K_C(T, P) \equiv \exp\left[-\frac{\sum_i v_i \mu_i^*}{k_B T}\right] \rho_{\text{solvent}}^{\sum_i v_i}$$

$$\left(\frac{\partial \ln K_x}{\partial(1/T)}\right)_P = -\frac{\Delta h^*}{k_B} \quad \text{where} \quad \Delta h^* \equiv \sum_i v_i h_i^* \quad \left(\frac{\partial \ln K_x}{\partial P}\right)_T = -\frac{\Delta v^*}{k_B T} \quad \text{where} \quad \Delta v^* \equiv \sum_i v_i v_i^*$$

Classical transition state theory:

$$\frac{\text{molecules reacted}}{\text{time} \cdot \text{volume}} = k_f \rho_A \rho_B \rho_C \quad \text{where} \quad k_f \equiv \left(\frac{k_B T}{h}\right) e^{-\beta \Delta F_f^\ddagger}$$

$$\Delta F_f^\ddagger \equiv -k_B T \ln \left[ \frac{(q^\ddagger/V)}{(q_A/V)(q_B/V)(q_C/V)} \right] + U^\ddagger$$

## Statistical mechanical ensembles

property	microcanonical	canonical	grand canonical	isothermal-isobaric
<b>constant conditions</b>	$E, V, N$	$T, V, N$	$T, V, \mu$	$T, P, N$
<b>fluctuations</b>	none	$E$	$E, N$	$E, V$
<b>microstate probabilities</b>	$\wp_m = \frac{\delta_{E_m, E}}{\Omega(E, V, N)}$	$\wp_m = \frac{e^{-\beta E_m}}{Q(T, V, N)}$	$\wp_m = \frac{e^{-\beta E_m + \beta \mu N_m}}{\Xi(T, V, \mu)}$	$\wp_m = \frac{e^{-\beta E_m - \beta P V_m}}{\Delta(T, P, N)}$
<b>partition function</b>	$\Omega(E, V, N) = \sum_n \delta_{E_n, E}$	$Q(T, V, N) = \sum_n e^{-\beta E_n}$	$\Xi(T, V, \mu) = \sum_N \sum_n e^{-\beta E_n + \beta \mu N}$	$\Delta(T, P, N) = \sum_V \sum_n e^{-\beta E_n - \beta P V}$
<b>relations to other partition functions</b>	---	$Q = \sum_E e^{-\beta E} \Omega$	$\Xi = \sum_N \lambda^N Q$ $= \sum_N \sum_E \lambda^N e^{-\beta E} \Omega$ $\lambda \equiv \exp[\beta \mu]$	$\Delta = \sum_V e^{-\beta P V} Q$ $= \sum_V \sum_E e^{-\beta E - \beta P V} \Omega$
<b>thermodynamic potential</b>	$S = k_B \ln \Omega(E, V, N)$	$A = -k_B T \ln Q(T, V, N)$	$PV = k_B T \ln \Xi(T, V, \mu)$	$G = -k_B T \ln \Delta(T, P, N)$
<b>classical partition function</b>	$\Omega = \frac{1}{h^{3N} N!} \int \delta[H(\mathbf{p}^N, \mathbf{r}^N) - E] d\mathbf{p}^N d\mathbf{r}^N$	$Q = \frac{Z(T, V, N)}{\Lambda(T)^{3N} N!}$ $Z \equiv \int e^{-\beta U(\mathbf{r}^N)} d\mathbf{r}^N$ $\Lambda \equiv (h^2 / 2\pi m k_B T)^{\frac{1}{2}}$	$\Xi = \sum_{N=0}^{\infty} \frac{\lambda^N Z(T, V, N)}{\Lambda(T)^{3N} N!}$ $\lambda \equiv \exp[\beta \mu]$	$\Delta = \frac{1}{\Lambda(T)^{3N} N!} \int_0^{\infty} e^{-\beta P V} Z(T, V, N) dV$

\*Sums over  $n$  correspond to sums over all microstates at a given  $V$  and  $N$ .

\*\*Sums over  $N$  are from 0 to  $\infty$ , for  $V$  from 0 to  $\infty$ , and for  $E$  from  $-\infty$  to  $\infty$ .

\*\*\*Classical partition functions are given for a monatomic system of indistinguishable, structureless particles.